

Terminal singularities and $U(1)$ factors in F-theory

Andrew P. Turner
University of Pennsylvania

String Phenomenology '22

July 7, 2022

Based on: 22nn.nnnnn with A. Grassi, N. Raghuram, T. Weigand
2110.10159 with N. Raghuram

Introduction

- Singularities appear frequently in the geometric spaces relevant for physics, and so it behooves us to understand them from both a mathematical and physical perspective

Introduction

- Singularities appear frequently in the geometric spaces relevant for physics, and so it behooves us to understand them from both a mathematical and physical perspective
- Non-crepantly resolvable singularities in F-theory indicate the presence of localized uncharged massless matter states [Braun, Morrison '14] [Braun, Collinucci, Valandro '14] [Morrison, Park, Taylor '16] [Arras, Grassi, Weigand '16]

Introduction

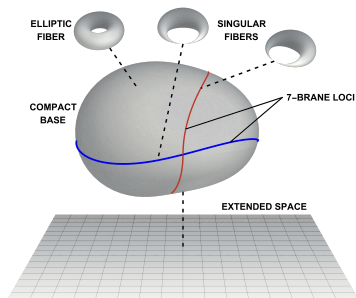
- Singularities appear frequently in the geometric spaces relevant for physics, and so it behooves us to understand them from both a mathematical and physical perspective
- Non-crepantly resolvable singularities in F-theory indicate the presence of localized uncharged massless matter states [Braun, Morrison '14] [Braun, Collinucci, Valandro '14] [Morrison, Park, Taylor '16] [Arras, Grassi, Weigand '16]
- Goal: investigate what happens when we introduce abelian gauge factors to F-theory models with non-crepant singularities

Introduction

- Singularities appear frequently in the geometric spaces relevant for physics, and so it behooves us to understand them from both a mathematical and physical perspective
- Non-crepantly resolvable singularities in F-theory indicate the presence of localized uncharged massless matter states [Braun, Morrison '14] [Braun, Collinucci, Valandro '14] [Morrison, Park, Taylor '16] [Arras, Grassi, Weigand '16]
- Goal: investigate what happens when we introduce abelian gauge factors to F-theory models with non-crepant singularities
- Brief advertisement of recent proposal of heuristic method to read off $U(1)$ charges without carrying out resolution, similar to Katz–Vafa

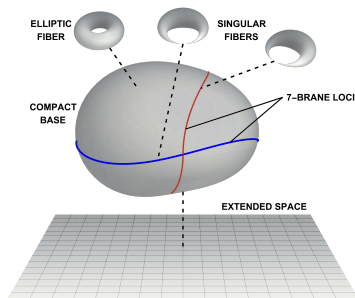
F-theory overview

- Elliptically fibered Calabi–Yau n -fold X :
 - ▶ Torus over each point in base B , $\pi: X \rightarrow B$
 - ▶ Has a section, $\sigma: B \rightarrow X$ s.t. $\pi \circ \sigma = \text{Id}_B$
 - ▶ Complex structure τ encodes Type IIB axiodilaton



F-theory overview

- Elliptically fibered Calabi–Yau n -fold X :
 - ▶ Torus over each point in base B , $\pi: X \rightarrow B$
 - ▶ Has a section, $\sigma: B \rightarrow X$ s.t. $\pi \circ \sigma = \text{Id}_B$
 - ▶ Complex structure τ encodes Type IIB axiodilaton



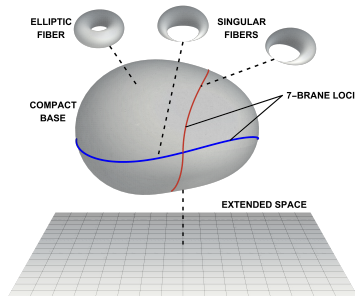
- Described by Weierstrass model: hypersurface

$$y^2 = x^3 + fxz^4 + gz^6$$

in ambient $\mathbb{P}_{[x:y:z]}^{2,3,1}$, where f, g are sections of $-4K_B, -6K_B$

F-theory overview

- Elliptically fibered Calabi–Yau n -fold X :
 - ▶ Torus over each point in base B , $\pi: X \rightarrow B$
 - ▶ Has a section, $\sigma: B \rightarrow X$ s.t. $\pi \circ \sigma = \text{Id}_B$
 - ▶ Complex structure τ encodes Type IIB axiodilaton



- Described by Weierstrass model: hypersurface

$$y^2 = x^3 + fxz^4 + gz^6$$

in ambient $\mathbb{P}_{[x:y:z]}^{2,3,1}$, where f, g are sections of $-4K_B, -6K_B$

- Gauge algebras and matter:
 - ▶ Codimension-one singularities (7-branes) \longleftrightarrow nonabelian gauge algebras
 - ▶ Additional rational sections $\longleftrightarrow \mathfrak{u}(1)$ gauge algebras
 - ▶ Codimension-two singularities \longleftrightarrow massless matter

Singularities

- A singular CY X can always be resolved: $\rho: \tilde{X} \rightarrow X$ such that \tilde{X} is smooth

Singularities

- A singular CY X can always be resolved: $\rho: \tilde{X} \rightarrow X$ such that \tilde{X} is smooth
- \tilde{X} and X agree on dense open sets, but not on the exceptional locus
 - ▶ Big resolution: exceptional locus contains divisors E_i
 - ▶ Small resolution: exceptional locus contains no divisors

Singularities

- A singular CY X can always be resolved: $\rho: \tilde{X} \rightarrow X$ such that \tilde{X} is smooth
- \tilde{X} and X agree on dense open sets, but not on the exceptional locus
 - ▶ Big resolution: exceptional locus contains divisors E_i
 - ▶ Small resolution: exceptional locus contains no divisors
- $K_{\tilde{X}} = K_X + \sum_i a_i E_i$, a_i “discrepancies”
 - ▶ $a_i \geq 0$: canonical singularity
 - ▶ $a_i > 0$: terminal singularity
 - ▶ $a_i = 0$: crepantly resolvable $\Rightarrow \tilde{X}$ is CY

Singularities

- A singular CY X can always be resolved: $\rho: \tilde{X} \rightarrow X$ such that \tilde{X} is smooth
- \tilde{X} and X agree on dense open sets, but not on the exceptional locus
 - ▶ Big resolution: exceptional locus contains divisors E_i
 - ▶ Small resolution: exceptional locus contains no divisors
- $K_{\tilde{X}} = K_X + \sum_i a_i E_i$, a_i “discrepancies”
 - ▶ $a_i \geq 0$: canonical singularity
 - ▶ $a_i > 0$: terminal singularity
 - ▶ $a_i = 0$: crepantly resolvable $\Rightarrow \tilde{X}$ is CY
- F-theory: resolution of singular point in fiber introduces $\text{rk } g$ exceptional \mathbb{P}^1 s
 - ▶ Gauge fields: $C_3 = A^i \wedge [E_i]$ and wrapped M2-branes
 - ▶ Localized matter: M2-branes wrapped on additional fibral curves C at codimension two
 - ▶ Mass along dual M-theory Coulomb branch: $m_0 \sim \text{vol}(C)$

Non-crepant singularities

- All non-crepantly resolvable singularities in an elliptically fibered CY n -fold occur in codimension-two or higher

Non-crepant singularities

- All non-crepantly resolvable singularities in an elliptically fibered CY n -fold occur in codimension-two or higher
- In particular, we focus on isolated \mathbb{Q} -factorial terminal singularities in codimension two of elliptic CY threefolds X_3
 - ▶ Every Weil divisor is also \mathbb{Q} -Cartier, i.e., there exists some $r \in \mathbb{Z}$ so that rD is Cartier
 - ▶ No small resolution as well as no crepant big resolution

Non-crepant singularities

- All non-crepantly resolvable singularities in an elliptically fibered CY n -fold occur in codimension-two or higher
- In particular, we focus on isolated \mathbb{Q} -factorial terminal singularities in codimension two of elliptic CY threefolds X_3
 - ▶ Every Weil divisor is also \mathbb{Q} -Cartier, i.e., there exists some $r \in \mathbb{Z}$ so that rD is Cartier
 - ▶ No small resolution as well as no crepant big resolution
- Physical interpretation: KK zero mode remains massless everywhere in the M-theory Coulomb branch \Rightarrow localized uncharged matter in codimension two [Arras, Grassi, Weigand '16]

Non-crepant singularities

- All non-crepantly resolvable singularities in an elliptically fibered CY n -fold occur in codimension-two or higher
- In particular, we focus on isolated \mathbb{Q} -factorial terminal singularities in codimension two of elliptic CY threefolds X_3
 - ▶ Every Weil divisor is also \mathbb{Q} -Cartier, i.e., there exists some $r \in \mathbb{Z}$ so that rD is Cartier
 - ▶ No small resolution as well as no crepant big resolution
- Physical interpretation: KK zero mode remains massless everywhere in the M-theory Coulomb branch \Rightarrow localized uncharged matter in codimension two [Arras, Grassi, Weigand '16]
- Number of uncharged massless hypermultiplets counted by Milnor number m_P :

$$\text{CxDef}(X_3) = \underbrace{\left(\frac{1}{2}b_3(X_3) - 1 - \frac{1}{2} \sum_P m_P \right)}_{\text{def. independent of sing.}} + \underbrace{\sum_P m_P}_{\text{def. of sing.}}$$

Examples

Many examples discussed in literature:

Examples

Many examples discussed in literature:

- $I_1 \rightarrow I_2$ [Grassi, Morrison '00] [Braun, Collinucci, Valandro '14] [Martucci, Weigand '15]

Examples

Many examples discussed in literature:

- $I_1 \rightarrow I_2$ [Grassi, Morrison '00] [Braun, Collinucci, Valandro '14] [Martucci, Weigand '15]
- \mathbb{Z}_k Weierstrass models [Braun, Morrison '14] [Morrison, Taylor '14] [Anderson, Etxebarria, Grimm, Keitel '14] [Klevers, Mayorga Peña, Oehlmann, Piragua, Reuter '14] [Mayrhofer, Palti, Till, Weigand '14] [Cvetič, Donagi, Klevers, Piragua, Poretschkin '15] [Cvetič, Grassi, Poretschkin '16] [Braun, Cvetič, Dongai, Poretschkin '17] ...

Examples

Many examples discussed in literature:

- $I_1 \rightarrow I_2$ [Grassi, Morrison '00] [Braun, Collinucci, Valandro '14] [Martucci, Weigand '15]
- \mathbb{Z}_k Weierstrass models [Braun, Morrison '14] [Morrison, Taylor '14] [Anderson, Etxebarria, Grimm, Keitel '14] [Klevers, Mayorga Peña, Oehlmann, Piragua, Reuter '14] [Mayrhofer, Palti, Till, Weigand '14] [Cvetič, Donagi, Klevers, Piragua, Poretschkin '15] [Cvetič, Grassi, Poretschkin '16] [Braun, Cvetič, Dongai, Poretschkin '17] ...
- $II \rightarrow III$ or IV [Arras, Grassi, Weigand '16]

Examples

Many examples discussed in literature:

- $I_1 \rightarrow I_2$ [Grassi, Morrison '00] [Braun, Collinucci, Valandro '14] [Martucci, Weigand '15]
- \mathbb{Z}_k Weierstrass models [Braun, Morrison '14] [Morrison, Taylor '14] [Anderson, Etxebarria, Grimm, Keitel '14] [Klevers, Mayorga Peña, Oehlmann, Piragua, Reuter '14] [Mayrhofer, Palti, Till, Weigand '14] [Cvetič, Donagi, Klevers, Piragua, Poretschkin '15] [Cvetič, Grassi, Poretschkin '16] [Braun, Cvetič, Dongai, Poretschkin '17] ...
- $II \rightarrow III$ or IV [Arras, Grassi, Weigand '16]
- $\mathfrak{su}(2)$ $III \rightarrow I_0^*$ family [Arras, Grassi, Weigand '16]

Examples

Many examples discussed in literature:

- $I_1 \rightarrow I_2$ [Grassi, Morrison '00] [Braun, Collinucci, Valandro '14] [Martucci, Weigand '15]
- \mathbb{Z}_k Weierstrass models [Braun, Morrison '14] [Morrison, Taylor '14] [Anderson, Etxebarria, Grimm, Keitel '14] [Klevers, Mayorga Peña, Oehlmann, Piragua, Reuter '14] [Mayrhofer, Palti, Till, Weigand '14] [Cvetič, Donagi, Klevers, Piragua, Poretschkin '15] [Cvetič, Grassi, Poretschkin '16] [Braun, Cvetič, Dongai, Poretschkin '17] ...
- $II \rightarrow III$ or IV [Arras, Grassi, Weigand '16]
- $\mathfrak{su}(2)$ $III \rightarrow I_0^*$ family [Arras, Grassi, Weigand '16]
- $\mathfrak{sp}(k)$ $I_{2k+1}^{\text{nonsplit}} \rightarrow I_{2k+2}$ family [Grassi, Weigand '18]

Example [Arras, Grassi, Weigand '16]

Type III model with $\text{ord}(g) = 4$

Example [Arras, Grassi, Weigand '16]

Type III model with $\text{ord}(g) = 4$

- Gauge algebra $\mathfrak{g} = \mathfrak{su}(2)$

$$y^2 + a_{1,2}\sigma^2xyz + a_{3,2}\sigma^2yz^3 = x^3 + a_{2,3}\sigma^3x^2z^2 + a_{4,1}\sigma xz^4 + a_{6,4}\sigma^4z^6$$

Example [Arras, Grassi, Weigand '16]

Type III model with $\text{ord}(g) = 4$

- Gauge algebra $\mathfrak{g} = \mathfrak{su}(2)$

$$y^2 + a_{1,2}\sigma^2xyz + a_{3,2}\sigma^2yz^3 = x^3 + a_{2,3}\sigma^3x^2z^2 + a_{4,1}\sigma xz^4 + a_{6,4}\sigma^4z^6$$

- \mathbb{Q} -factorial terminal I_0^* at $\{\sigma = a_{4,1} = 0\}$ supporting $2 \times \mathbf{2} + \mathbf{1}$

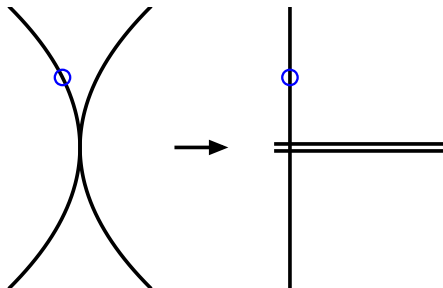
Example [Arras, Grassi, Weigand '16]

Type III model with $\text{ord}(g) = 4$

- Gauge algebra $\mathfrak{g} = \mathfrak{su}(2)$

$$y^2 + a_{1,2}\sigma^2xyz + a_{3,2}\sigma^2yz^3 = x^3 + a_{2,3}\sigma^3x^2z^2 + a_{4,1}\sigma xz^4 + a_{6,4}\sigma^4z^6$$

- \mathbb{Q} -factorial terminal I_0^* at $\{\sigma = a_{4,1} = 0\}$ supporting $2 \times \mathbf{2} + \mathbf{1}$



Example

Type III model with $\text{ord}(g) = 4$ and $\mathfrak{u}(1)$ factor

Example

Type III model with $\text{ord}(g) = 4$ and $\mathfrak{u}(1)$ factor

- Gauge algebra $\mathfrak{g} = \mathfrak{su}(2) \oplus \mathfrak{u}(1)$ [Küntzler, Schäfer-Nameki '14]

$$\begin{aligned} y^2 s + \tilde{b}_{0,1} \sigma x^2 y + \tilde{b}_{1,2} \sigma^2 w x y s + \tilde{b}_{2,2} \sigma^2 w^2 y s^2 \\ = \tilde{c}_{0,4} \sigma^4 w^4 s^3 + \tilde{c}_{1,1} \sigma w^3 x s^2 + \tilde{c}_{2,3} \sigma^3 w^2 x^2 s + \tilde{c}_3 w x^3 \end{aligned}$$

Example

Type III model with $\text{ord}(g) = 4$ and $\mathfrak{u}(1)$ factor

- Gauge algebra $\mathfrak{g} = \mathfrak{su}(2) \oplus \mathfrak{u}(1)$ [Küntzler, Schäfer-Nameki '14]

$$\begin{aligned} y^2 s + \tilde{b}_{0,1} \sigma x^2 y + \tilde{b}_{1,2} \sigma^2 w x y s + \tilde{b}_{2,2} \sigma^2 w^2 y s^2 \\ = \tilde{c}_{0,4} \sigma^4 w^4 s^3 + \tilde{c}_{1,1} \sigma w^3 x s^2 + \tilde{c}_{2,3} \sigma^3 w^2 x^2 s + \tilde{c}_3 w x^3 \end{aligned}$$

- \mathbb{Q} -factorial terminal I_0^* at $\{\sigma = \tilde{c}_{1,1} = 0\}$ supporting $2 \times \mathbf{2}_0 + \mathbf{1}_0$

Example

Type III model with $\text{ord}(g) = 4$ and $\mathfrak{u}(1)$ factor

- Gauge algebra $\mathfrak{g} = \mathfrak{su}(2) \oplus \mathfrak{u}(1)$ [Küntzler, Schäfer-Nameki '14]

$$\begin{aligned} y^2 s + \tilde{b}_{0,1} \sigma x^2 y + \tilde{b}_{1,2} \sigma^2 w x y s + \tilde{b}_{2,2} \sigma^2 w^2 y s^2 \\ = \tilde{c}_{0,4} \sigma^4 w^4 s^3 + \tilde{c}_{1,1} \sigma w^3 x s^2 + \tilde{c}_{2,3} \sigma^3 w^2 x^2 s + \tilde{c}_3 w x^3 \end{aligned}$$

- \mathbb{Q} -factorial terminal I_0^* at $\{\sigma = \tilde{c}_{1,1} = 0\}$ supporting $2 \times \mathbf{2}_0 + \mathbf{1}_0$
- Crepantly resolvable I_0^* at $\{\sigma = \tilde{c}_3 = 0\}$ supporting $2 \times \mathbf{2}_1 + \mathbf{1}_2$

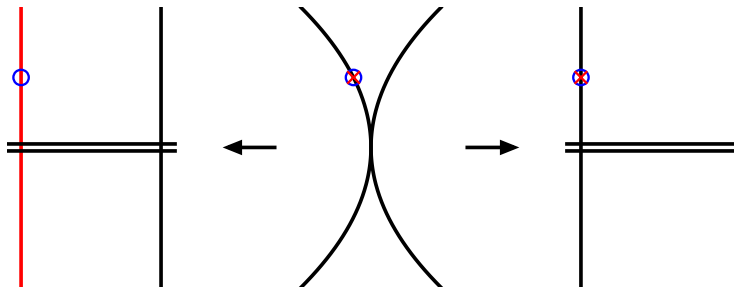
Example

Type III model with $\text{ord}(g) = 4$ and $\mathfrak{u}(1)$ factor

- Gauge algebra $\mathfrak{g} = \mathfrak{su}(2) \oplus \mathfrak{u}(1)$ [Küntzler, Schäfer-Nameki '14]

$$\begin{aligned} y^2 s + \tilde{b}_{0,1} \sigma x^2 y + \tilde{b}_{1,2} \sigma^2 w x y s + \tilde{b}_{2,2} \sigma^2 w^2 y s^2 \\ = \tilde{c}_{0,4} \sigma^4 w^4 s^3 + \tilde{c}_{1,1} \sigma w^3 x s^2 + \tilde{c}_{2,3} \sigma^3 w^2 x^2 s + \tilde{c}_3 w x^3 \end{aligned}$$

- \mathbb{Q} -factorial terminal I_0^* at $\{\sigma = \tilde{c}_{1,1} = 0\}$ supporting $2 \times \mathbf{2}_0 + \mathbf{1}_0$
- Crepantly resolvable I_0^* at $\{\sigma = \tilde{c}_3 = 0\}$ supporting $2 \times \mathbf{2}_1 + \mathbf{1}_2$



Reading off $U(1)$ charges [Raghuram, APT '21]

Consider a model with gauge algebra $\mathfrak{g} \oplus \mathfrak{u}(1)$, with \mathfrak{g} simply laced

- Codimension-one: singularity types G_i at loci $\{\sigma_i = 0\}$
- Codimension-two: singularity type enhances $G_{i_1} \times \cdots \times G_{i_N} \rightarrow H$

Reading off $U(1)$ charges [Raghuram, APT '21]

Consider a model with gauge algebra $\mathfrak{g} \oplus \mathfrak{u}(1)$, with \mathfrak{g} simply laced

- Codimension-one: singularity types G_i at loci $\{\sigma_i = 0\}$
- Codimension-two: singularity type enhances $G_{i_1} \times \cdots \times G_{i_N} \rightarrow H$

Codimension-one orders of vanishing:

$$\text{ord}_1(\hat{z}) = 0, \quad (\text{ord}_1(\hat{x}), \text{ord}_1(\hat{y}), \text{ord}_1(3\hat{x}^2 + f\hat{z}^4)) = \vec{\tau}_G(\nu)$$

Codimension-two orders of vanishing:

$$\text{ord}_2(\hat{z}) = \frac{1}{2} \left(\frac{\prod_i d_{G_i}}{d_H} q^2 + \sum_i T_{G_i}(\nu_i) - T_H(\mu) \right)$$

$$(\text{ord}_2(\hat{x}), \text{ord}_2(\hat{y}), \text{ord}_2(3\hat{x}^2 + f\hat{z}^4)) = (2, 3, 4) \times \text{ord}_2(\hat{z}) + \vec{\tau}_H(\mu)$$

q : $\mathfrak{u}(1)$ charge

$[\hat{x} : \hat{y} : \hat{z}]$: section components

ν_i, μ : Integers denoting the components of the resolved fiber hit by the generating section

$\vec{\tau}_G(\nu)$: Triplet of integers dependent on G and ν

d_G : Number of elements in the center of G

$T_G(\nu) = (C_G^{-1})_{\nu\nu}$: ν th diagonal entry of inverse Cartan matrix of G

Conclusions

- The introduction of a $U(1)$ factor to a model with \mathbb{Q} -factorial terminal singularities can give charge to previously uncharged localized massless hypermultiplets, changing the nature of the singularity to allow for a crepant resolution

Conclusions

- The introduction of a $U(1)$ factor to a model with \mathbb{Q} -factorial terminal singularities can give charge to previously uncharged localized massless hypermultiplets, changing the nature of the singularity to allow for a crepant resolution
- Can this be interpreted as an “enhancement” of discrete symmetries (possibly \mathbb{Z}_1) under which the localized multiplets are charged?

Conclusions

- The introduction of a $U(1)$ factor to a model with \mathbb{Q} -factorial terminal singularities can give charge to previously uncharged localized massless hypermultiplets, changing the nature of the singularity to allow for a crepant resolution
- Can this be interpreted as an “enhancement” of discrete symmetries (possibly \mathbb{Z}_1) under which the localized multiplets are charged?
- What about models that are forced to give nonzero $U(1)$ charge to all nonabelian-charged hypermultiplets?

Euler characteristic

$$\frac{1}{2}\chi_{\text{top}} = \text{KaDef}(X) - \text{CxDef}(X) + \frac{1}{2}m_P$$

Compute:

- Resolve all codimension-one singularities
- Sum over all fibers:

$$\begin{aligned}\chi_{\text{top}}(X_3) = & \sum_i \chi_{\text{top}}(X_{P_i}) \cdot B_i + \chi_{\text{top}}(X_{\Sigma_1}) \cdot (2 - 2g(\Sigma_1) - \sum_i B_i) \\ & + 2 - 2g(\Sigma_0) + 3C + \sum_i \epsilon_i B_i\end{aligned}$$

with

$$C = (-4K_B - \text{ord}(f))(-6K_B - \text{ord}(g)) - \sum_i \mu_{P_i}(f, g) B_i$$